Proper Heavy $Q\bar{Q}$ Potential
from Lattice QCD

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See also: A.R., T. Hatsuda, S. Sasaki arXiv:0910.2321

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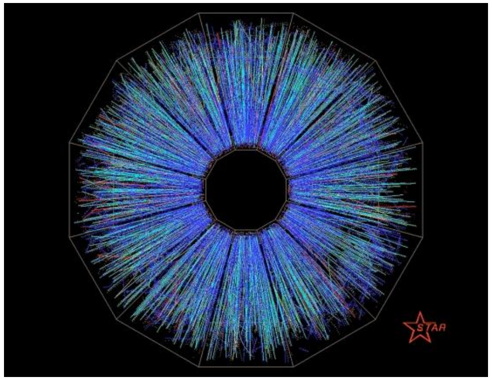
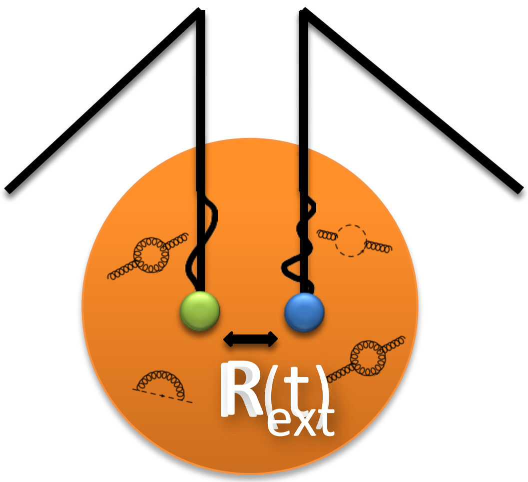
Research Interests:

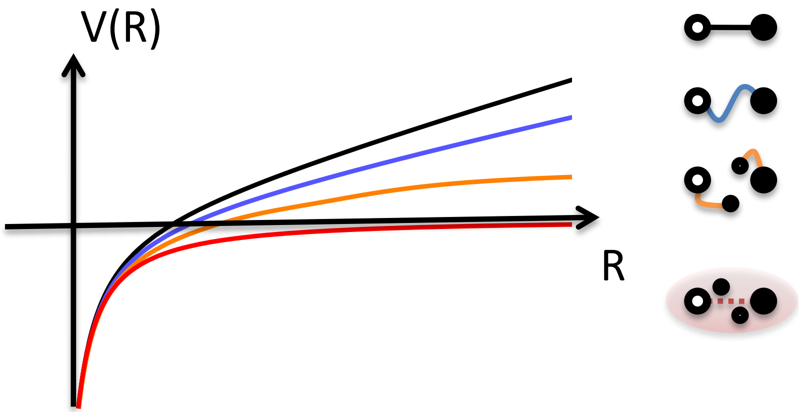
- Heavy Ion Collisions: The Quark Gluon Plasma
  - **QGP**: High temperature state where the fundamental constituents of matter, i.e. quarks and gluons are not confined within Hadrons
  - **Goal**: Understanding from 1st principles the properties of heavy quark bound states (J/Ψ) as they cross into the QGP phase
  - **Tools**: Lattice QCD simulations, i.e Monte Carlo Integration of spatially regularized SU(3) fields



BNL,RHIC, Star Experiment

Proper Heavy $Q\bar{Q}$ Potential
from Lattice QCD





Heavy Quarkonia @ finite T

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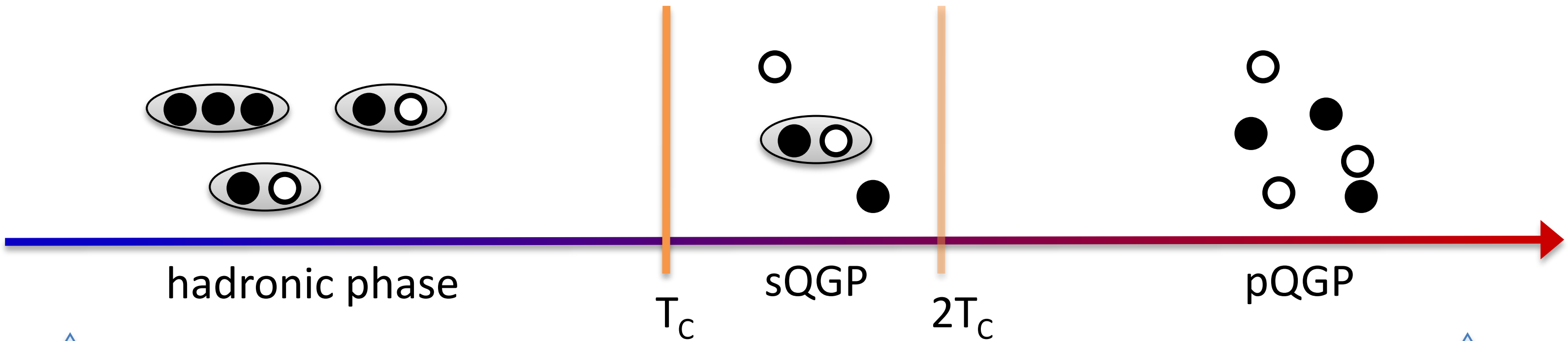
- m ≫ T : non-relativistic description using a static potential



T=0 case Wilson loop	Potential Models	Hard Thermal Loop
$\psi(t) = e^{imt}\, e^{ig\int_0^t A^0(t)dt}\psi(0)$	$V(R) \sim \Delta F_{Q\bar{Q}}(R) \in \mathbb{R}$	$V(R) = V_{DS}(R) + iV_{LD}(R)$
$V(R) \overset{\tau\to\infty}{=} -\frac{1}{\tau}\log\left[WL(\tau,R)\right]$	$= -\frac{1}{\beta}\log\left[\langle L(R)L^\dagger(0)\rangle\right]$	$Q\bar{Q}$ is transient state

We propose a GAUGE-INVARIANT and NON-PERTURBATIVE definition of the proper potential based on the spectral function

Outline

- Motivation: Heavy Quark Potential
- Proper Potential from Lattice QCD
  - Idea and formulation
  - First results for $Re[V_0(R)]$
- Conclusion and Outlook

Formulation

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- Starting point is the QQbar correlator and its spectral function:
  - QQbar mesic operator: $M(x,y) = \bar{\psi}(x)\,\Gamma\, U(x,y)\psi(y)$
  - QQbar forward correlator: $D^>(t,R) = \langle M_R(t) M_R^\dagger(0)\rangle$
  - Spectral function at finite T:

$$\rho(\omega,R) = \frac{1}{Z}\sum_{n,n'} |\langle n|M(0)|n'\rangle|^2 \left(e^{-\beta E_n} - e^{-\beta E_{n'}}\right)\delta(E_{n'} - E_n - \omega)$$

Antisymmetry

- In the spectral function we find three mutually exclusive cases:
  - Case I
  - Case II
  - Case III

Case I	Case II	Case III
Pure **medium** in $\vert n\rangle$	Pure **medium** in $\vert n'\rangle$	Q **or** Qbar +medium in $\vert n'\rangle$
QQbar + medium in $\vert n'\rangle$	**QQbar + medium** in $\vert n\rangle$	Qbar **or** Q +medium in $\vert n\rangle$

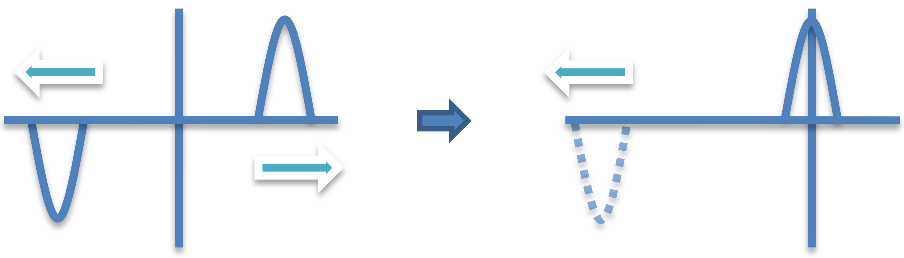
redundant

No interaction: irrelevant

Formulation II

- Preparations for a consistent $m \to \infty$ limit: frequency shift
  - Non-interacting case, two delta peaks
  - Retain positive ω peak: case I

$$\bar{\omega} = \omega - 2m$$



- Physics of the interaction lies in the relative position of the peak to $\bar{\omega}=0$

$$\rho^I(\bar{\omega}, R) = \frac{1}{Z_0} \sum_{n,n'} |\langle n|M(0)|Q\bar{Q} \in n'\rangle|^2 \delta(\bar{\omega} - (\epsilon_{n'}(R) - \epsilon_n)) e^{-\beta E_n}$$

  - Finite Temperature effects from sum over n: $\epsilon_{n'}(R) = E_{n'}(R) - 2m$ is T independent

- Heavy mass limit: retarded and forward correlator are equal

$$D^I_>(t) = \int_{-\infty}^{\infty} e^{-i\omega t} D^I_>(\omega) = \int_{-\infty}^{\infty} e^{-i\omega t} \rho^I(\omega) = e^{-2imt} \int_{-\infty}^{\infty} e^{-i\bar{\omega} t} \rho^I(\bar{\omega})$$

The Schrödinger Equation

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$$i\partial_t D^I_>(t) = e^{-2imt}\int_{-\infty}^{\infty} d\bar\omega \Big(2m+\bar\omega\Big) e^{-i\bar\omega t}\rho^I(\bar\omega)$$

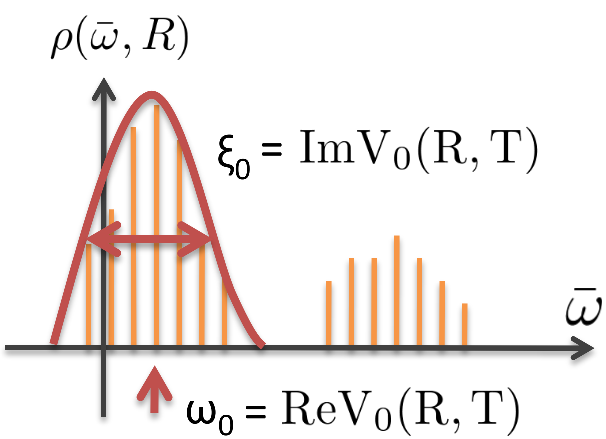
$$= 2mD^I_>(t) + e^{-2imt}\int_{-\infty}^{\infty} d\bar\omega \Big(\bar\omega\Big) e^{-i\bar\omega t}\rho^I(\bar\omega)$$

- Non interacting case:

$$\rho^I(\bar\omega) \propto \delta(\bar\omega)$$

$$i\partial_t D^I_>(t) = 2mD^I_>(t)$$

- In a finite volume all energies are discrete but their envelope can exhibit broad peaks



- „Ground state potential“: lowest lying peak structure
- Validity of Schrödinger description can be checked:

$$\int_{-\infty}^{\infty} d\bar\omega \Big(\bar\omega\Big) e^{-i\bar\omega t}\rho^I(\bar\omega, R) \stackrel{!}{=} V(R)\int_{-\infty}^{\infty} d\bar\omega e^{-i\bar\omega t}\rho^I(\bar\omega, R)$$

- In the high T region e.g. a Breit Wigner shape

$$\boxed{i\partial_t D^I_>(t) = \Big[2m + \mathrm{ReV}_0(R,T) - i\mathrm{ImV}_0(R,T)\Big] D^I_>(t)}$$

- Note: There is **no Schrödinger** equation for the **full** D>(t)

$$D_>(t) = e^{-i\omega_0 t} - e^{i\omega_0 t} \Rightarrow i\partial_t D_>(t) = \omega_0\Big[e^{-i\omega_0 t} + e^{i\omega_0 t}\Big]$$

Connection to Lattice QCD

- Analytic continuation gives:

$$D^I_>(\tau,R) = e^{-2m\tau}\int_{-\infty}^{\infty} d\bar\omega\, e^{-\bar\omega\tau}\rho^I(\bar\omega,R)$$

  - Cannot yet take the heavy mass limit
  - What form does $D^I_>$ have?

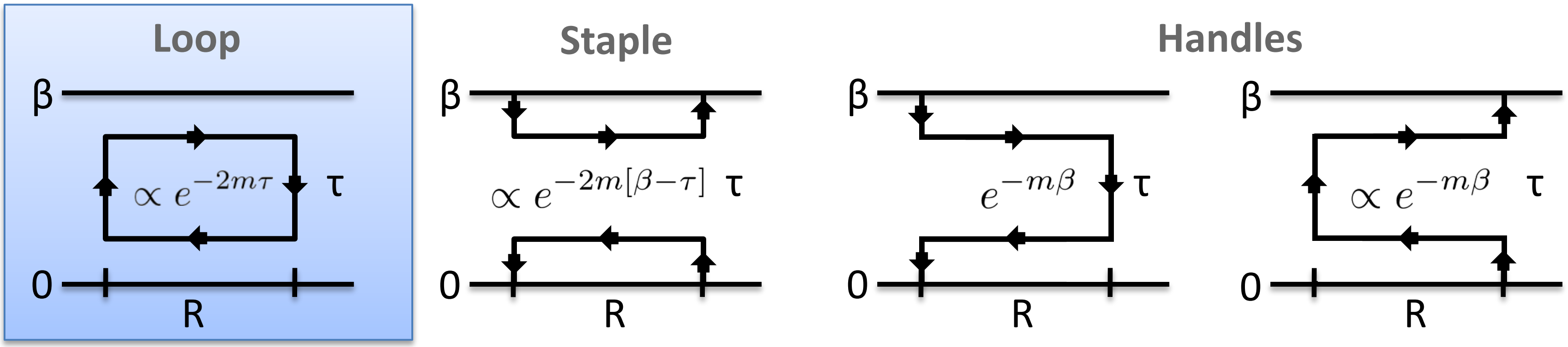
- Finding an expression for $D_>(\tau)$:

$$\langle M_R(\tau)M_R^\dagger(0)\rangle = -\sum_n \langle n|e^{-\beta H}\psi(x,0)\bar\psi(x,\tau)\Gamma U_\tau(\vec x,\vec y)\psi(y,\tau)\bar\psi(y,0)U_0(\vec y,\vec x)\bar\Gamma|n\rangle$$

$$S_E(z,z') = \langle T_\tau\,\psi(z,\tau)\bar\psi(z',\tau')\rangle \xrightarrow{\text{static limit}} (-i\gamma_4 D_\tau + m)S_E(\tau,\tau') = \delta(\tau-\tau')$$

  - Additionally: boundary conditions in τ direction

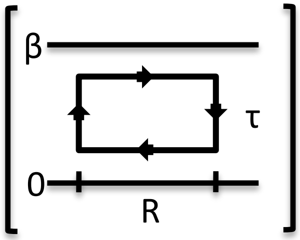
- In the **heavy mass limit**: only **Wilson lines** remain for $D^>(\tau)$



Exploring the Proper Potential

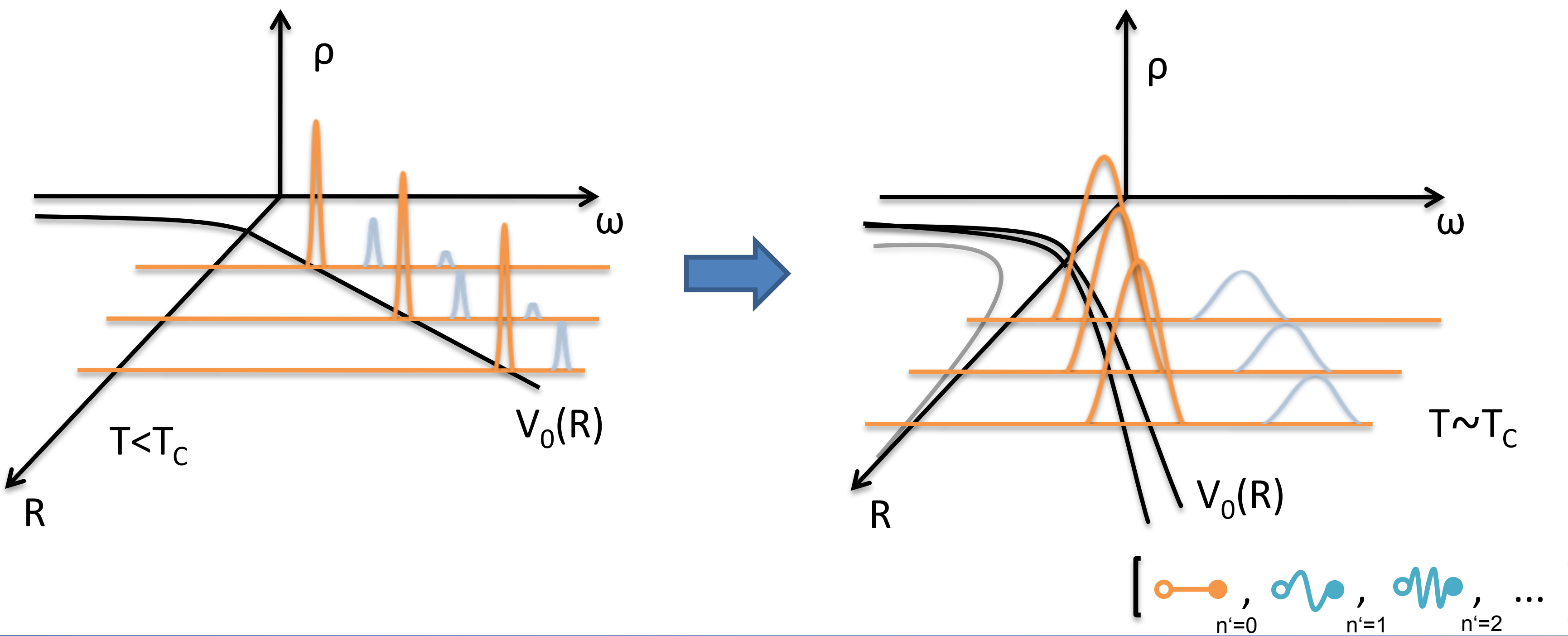
- Combining both spectral function and the imaginary time correlator:

$$\frac{D^{\text{loop}}_{>}(\tau,R)}{e^{-2m\tau}} = \text{Tr}\left[\ \cdots\ \right] = \int_{-\infty}^{\infty} d\bar{\omega}\, e^{-\tau\bar{\omega}} \rho^{\text{loop}}(\bar{\omega},R)$$



Lattice QCD — **MEM / Exponential Fitting**

- **Reconstruct** spectral function for **different R**: map the shape of the potential:

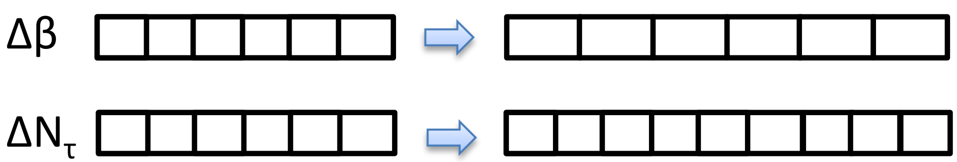


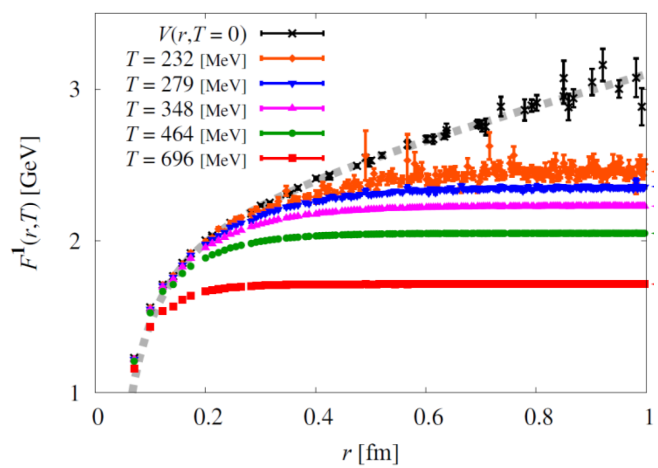
Pure SU(3): naive Wilson action

NX=20 NT=32 $T=2.33T_C$
NT=96 $T=0.78T_C$

$\beta=7.0$ $\xi_0=3.5$ $a_\tau= ¼a_\sigma =0.01\text{fm}$

- Monte Carlo Simulation: CPS++ Library
  - Action is local since no dynamical fermions present
  - Pseudo Heatbath Algorithm (Cabibbo - Marinari): „bring the current link in contact with thermal bath“
  - Overrelaxation for better decorrelation: additional single Metropolis step is inserted
- Anisotropic Lattices
  - Since τ direction is compact: need to increase # of points
  - Renormalized anisotropy needs to be determined: spatial and temporal Wilson Loop ratios
- Fixed scale vs. Fixed size Lattice QCD
  - Change NT instead of a_τ : **$T = 1/(NT \times a_\tau)$**
  - Renormalization parameter same for all T

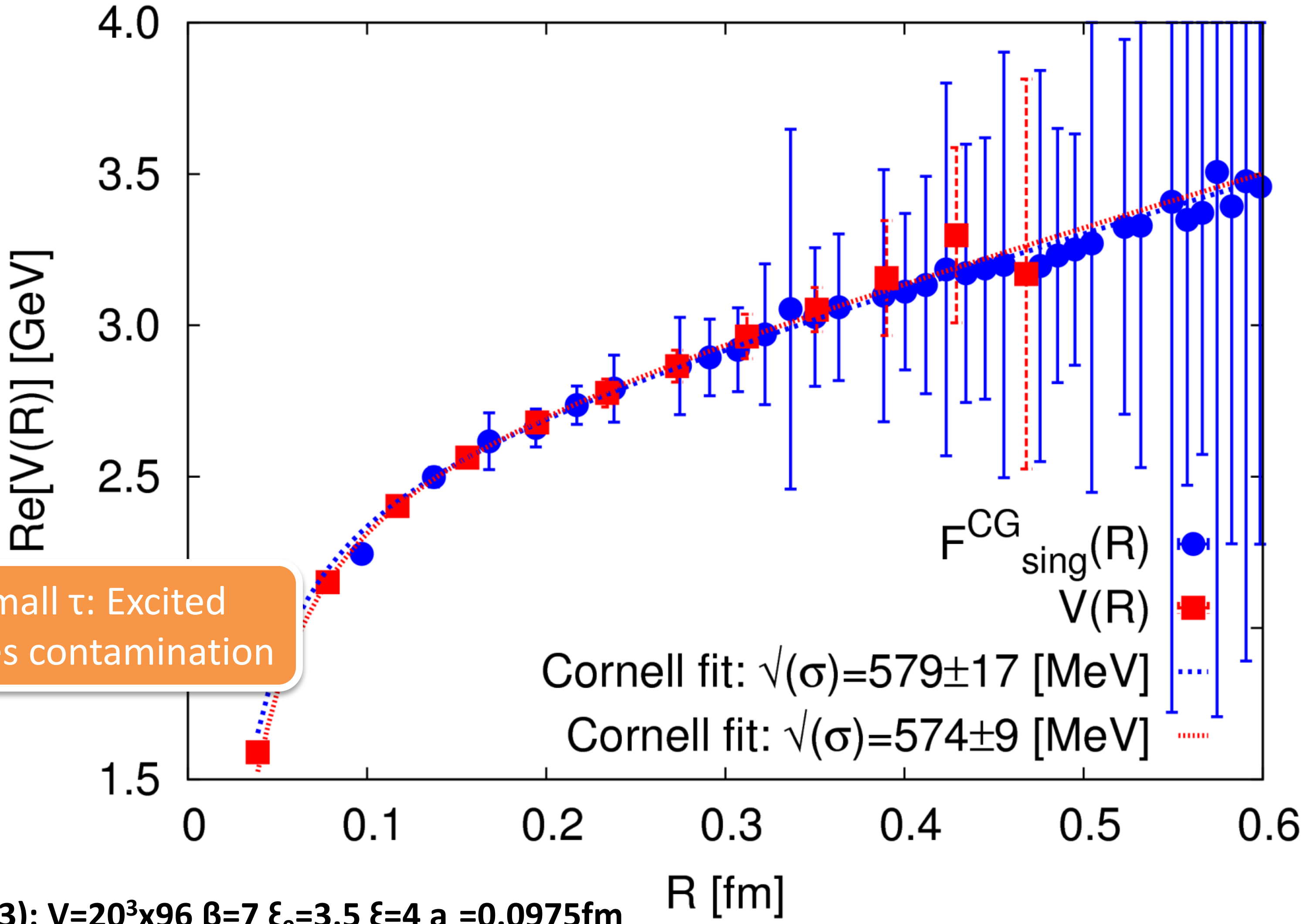




Maezawa et. al. PoS Lat 2009

small R physics is T independent

Simulation results at T=0.78T_c



Small τ: Excited states contamination

ential,
y part

Pure SU(3): V=20^3x96 β=7 ξ_0=3.5 ξ=4 a_τ=0.0975fm

Simulation results at $T=2.33T_C$

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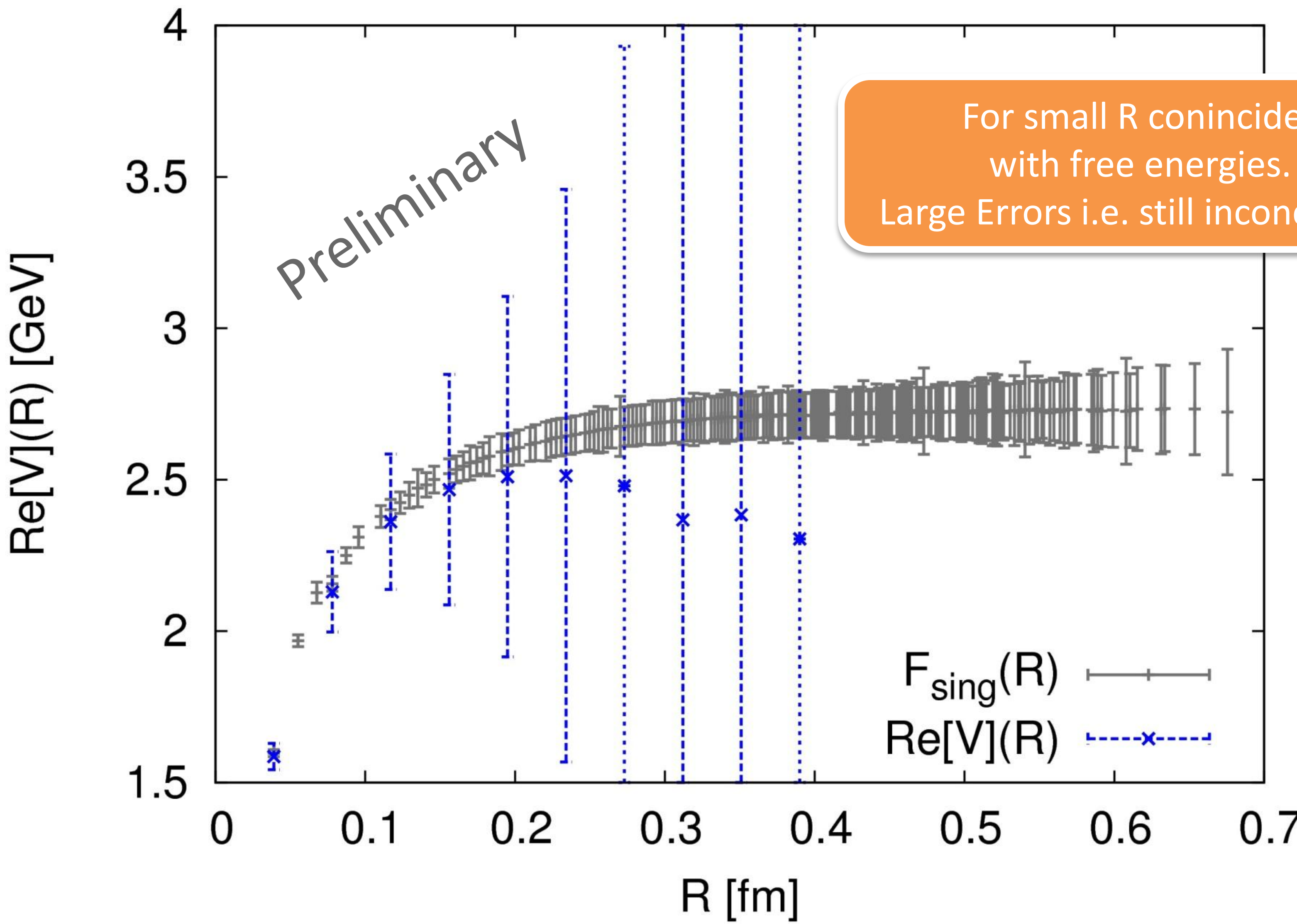

For small R conincides
with free energies.
Large Errors i.e. still inconclusive

ution
(ω̄)

Pure

Conclusion & Outlook

- First-principles definition of the heavy quark potential:
  - Wilson Loop and its spectral function connected to V(R,T) **non-perturbatively**
  - Ground state peak envelope of the spectral function leads to a **Schrödinger Equation**
  - Peak structure of the spectral function provides **real part** (position) and **imaginary part** (width) of the potential V(R,T)
  - Current results with $\mathbf{m_Q=\infty}$ and **quenched QCD:**

    Re[V(R,T)] below $\mathbf{T_C}$: Confining potential, , coincides with color singlet free energies

    Re[V(R,T)] above $\mathbf{T_C}$: Reconstruction of the proper potential not conclusive yet

- Work in progress and future directions:
  - Higher statistics data necessary: Using T2K for quenched QCD configuration generation
  - Full QCD simulations to include the influence of light quarks in the medium: High cost

The End

Thank you for your support and attention